

# Exterior Calculus

- Manifold - An  $n$ -dimensional (differentiable) manifold  $M^n$  is a space that is locally  $\mathbb{R}^n$  in the following sense; it is covered by a family of curvilinear coordinate systems  $\{U_\alpha; \varphi_\alpha\}$   $\varphi_\alpha: U_\alpha \rightarrow \mathbb{R}^n$  such that a point  $p \in U_\alpha \cap U_\beta$  that lies in two coordinate patches will have its two sets of coordinates related differentially

$$\varphi_\beta \circ \varphi_\alpha^{-1}: \mathbb{R}^n \rightarrow \mathbb{R}^n$$

$$x_\beta^i(p) = f_{\beta\alpha}^i(x_\alpha^1(p), \dots, x_\alpha^n(p)) \quad \text{for } i=1, \dots, n$$

(if the functions  $f_{\beta\alpha}^i$  are  $C^\infty$ , or real analytic we say that  $M$  is  $C^\infty$  or real analytic.)

- Vector - A tangent, or contravariant vector or vector at  $p_0 \in M^n$  call it  $v$  assigns to each coordinate system  $(U_\alpha, \varphi_\alpha)$  holding  $p_0$  an  $n$ -tuple  $(v_\alpha^i)$ ;  $v_\alpha^1, v_\alpha^2, \dots, v_\alpha^n$  such that if  $p_0 \in U_\alpha \cap U_\beta$

$$v_\beta^i = \sum_{j=1}^n \frac{\partial x_\beta^i}{\partial x_\alpha^j}(p_0) v_\alpha^j$$

associated with this vector is a differential operator

$D_v$  defined by

$$D_v(f) \equiv \sum_{i=1}^n \frac{\partial f}{\partial x^i}(p_0) v^i \equiv v_{p_0}(f)$$

It is easily verified that this is independent of the coordinate system.

In a specific coordinate system

$$v_{p_0} = \sum_{i=1}^n v^i \frac{\partial}{\partial x^i} \Big|_{p_0}$$

Tangent Space - to  $M^n$  at the point  $p \in M^n$  is the vector space  $M_p^n$  consisting of all tangent vectors to  $M^n$  at  $p$ . If  $(u, \varphi)$  is a coordinate system for  $M^n$  holding  $p$  the  $n$  vectors

$$e_1 = \left. \frac{\partial}{\partial x^1} \right|_p, \dots, e_n = \left. \frac{\partial}{\partial x^n} \right|_p$$

form a basis for  $M_p^n$  called a coordinate basis.  
 $M_p^n$  is an  $n$ -dimensional vector space.

Differential of a mapping  $M^n$  and  $V^r$  are manifolds  
 $F: M^n \rightarrow V^r$  is a map described in terms of local coordinates  $x$  near  $p \in M^n$  and  $y$  near  $F(p) \in V^r$  by  $n$  functions of  $n$  variables

$$y^d = F^d(x^1, \dots, x^n)$$

$F$  is a diffeomorphism provided  $F$  is 1:1 and onto and both  $F$  and  $F^{-1}$  are differentiable

The differential of the map  $F_*: M_p^n \rightarrow V_{F(p)}^r$  takes tangent vectors from  $M_p^n$  to tangent vectors in  $V_{F(p)}^r$  according to the following rule given the above local coordinate systems.

$$w^d = \sum_{i=1}^n \frac{\partial F^d}{\partial x^i}(x_0) v^i \quad F_*(v) = w$$

Definition of Dual space and covariant vector

Let  $E$  be a real vector space. If  $E$  is  $n$ -dimensional, let  $e_1, \dots, e_n$  be a basis. Thus if  $v \in E$   $v$  has a unique expansion

$$v = e_i v^i \quad (\text{Einstein summation convention})$$

$$\text{or } e = (e_1, \dots, e_n) \quad v = \begin{pmatrix} v^1 \\ \vdots \\ v^n \end{pmatrix} \Rightarrow v \equiv e v$$

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A real linear functional  $d$  on  $E$  is a linear transformation  
 $d: E \rightarrow \mathbb{R}$  ( $\mathbb{R}$ ,  $d(c_1V + c_2W) = c_1d(V) + c_2d(W)$   $V, W \in E$   $c_1, c_2 \in \mathbb{R}$ ),

If  $E$  is a vector space the collection of all linear functionals  $d$  on  $E$  defines a new vector space  $E^*$  the dual space to  $E$  under operations

$$(d + \beta)(V) = d(V) + \beta(V)$$

$$(cd)(V) = c d(V)$$

If  $e_1, \dots, e_n$  is a basis for  $E$  define the dual basis  $\sigma^1, \dots, \sigma^n$  of  $E^*$  by

$$\sigma^i(e_j) = \delta_j^i$$

$\Sigma$  note that linear functional  $\beta$  can be expanded in basis  
by  $\beta(V) = \beta(e_i \sigma^i) = \beta(e_i) \sigma^i(V) = \beta(e_i) \delta_j^i \sigma^j(V) = \beta(e_i) \sigma^i(e_j) \sigma^j(V)$   
 $= [\beta(e_i) \sigma^i](e_j \sigma^j)$

$$\Rightarrow \beta = \beta(e_i) \sigma^i = b_i \sigma^i \text{ where } b_i = \beta(e_i)$$

in above notation

$$b = (b_1, \dots, b_n) \quad \sigma = \begin{pmatrix} \sigma^1 \\ \vdots \\ \sigma^n \end{pmatrix} \Rightarrow \beta = b \sigma$$

$$\text{now } \beta(V) = b \sigma e V = b I V = b V \quad \}$$

If  $f: M^n \rightarrow \mathbb{R}$  its differential  $df$  is the linear functional on tangent spaces  $df: M_p^n \rightarrow \mathbb{R}$  given by

$$df(V_p) = V_p(f)$$

In coordinates:  $e_i = \frac{\partial}{\partial x^i} \Big|_p, \dots, e_n = \frac{\partial}{\partial x^n} \Big|_p$  form a basis for  $M_p^n$

$$\text{and } df(V_p) = V_p(f) = \left[ v^i(p) \frac{\partial}{\partial x^i} \Big|_p \right] (f) = v^i(p) \frac{\partial f}{\partial x^i} \Big|_p$$

$$\text{Let } f = x^i. \quad dx^i(e_j) = dx^i \left( \frac{\partial}{\partial x^j} \right) = \frac{\partial x^i}{\partial x^j} = \delta_j^i$$

$\therefore dx^1, \dots, dx^n$  form a dual basis to  $\frac{\partial}{\partial x^1}, \dots, \frac{\partial}{\partial x^n}$   
in other words  $dx^1, \dots, dx^n$  form a basis for  $(M_p^n)^*$   
the cotangent space

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Now for linear functional  $\alpha: M^n \rightarrow \mathbb{R}$  called a covariant vector or covector or 1-form can be expanded

$$\alpha = a_i dx^i \quad \text{where} \quad a_i \equiv \alpha\left(\frac{\partial}{\partial x^i}\right)$$

The differential is now naturally expanded

$$df = df\left(\frac{\partial}{\partial x^i}\right) dx^i = \frac{\partial f}{\partial x^i} dx^i$$

Under a change of coordinates  $d\bar{x}^i = d\bar{x}^j \left(\frac{\partial x^i}{\partial \bar{x}^j}\right) dx^j = \frac{\partial \bar{x}^i}{\partial x^j} dx^j$

For a general 1-form  $a_i dx^i = a_i \frac{\partial x^i}{\partial \bar{x}^j} d\bar{x}^j = \bar{a}_j d\bar{x}^j$

$\Rightarrow \bar{a}_j = \frac{\partial x^i}{\partial \bar{x}^j} a_i$  or in terms used for vectors

$$a_i^{\beta} = \sum_{j=1}^n a_j^{\alpha} \frac{\partial x_j^{\alpha}}{\partial x_i^{\beta}}$$

Pull-back of a mapping  $\alpha$ : for the differential of a mapping  $M^n$  and  $V^r$  are manifolds  $F: M^n \rightarrow V^r$  is a map described in terms of local coordinates  $x$  near  $p \in M^n$  and  $y$  near  $F(p) \in V^r$  by  $r$  functions of  $n$  variables

$$y^a = F^a(x^1, \dots, x^n)$$

The pull-back of the map  $F^*: (V_y^r)^* \rightarrow (M_{F^{-1}(y)}^n)^*$  takes covectors from  $(V_y^r)^*$  to covectors in  $(M_{F^{-1}(y)}^n)^*$  according to the following rule

$$(F^* \beta)(v) = \beta(F_* v) \quad \text{where} \quad v \in M_{F^{-1}(y)}^n \quad \beta \in (V_y^r)^* \\ F_* v \in V_y^r \quad F^* \beta \in (M_{F^{-1}(y)}^n)^*$$

in local coordinates

$$d_i = \frac{\partial F^a}{\partial x^i}(F^{-1}(y_0)) \beta_a \quad F^*(\beta) = d$$

in terms of basis vectors the transformations are

$$F^*(dy^a) = \frac{\partial y^a}{\partial x^i} dx^i \quad F_*\left(\frac{\partial}{\partial x^i}\right) = \frac{\partial y^a}{\partial x^i} \frac{\partial}{\partial y^a}$$

for the composition of two mappings  $F: M^n \rightarrow V^r$   $G: V^r \rightarrow U^p$

$$(G \circ F)^* = F^* \circ G^* \quad (G \circ F)_* = G_* \circ F_*$$

$$G \circ F: M^n \rightarrow U^p$$

• Tangent and Cotangent Bundles - the tangent bundle to a differentiable manifold  $M^n$  is the space of all tangent vectors to  $M^n$ . A point in this new space consists of a point  $(x, v)$  where  $x$  is a point of  $M^n$  and  $v$  is a tangent vector to  $M^n$  at  $x$  (i.e.,  $v \in M_x^n$ ). Denote the tangent bundle to  $M^n$  by  $TM^n$ .

We can introduce local coordinates in  $TM^n$  as follows. Let  $(x, v) \in TM^n$ .  $x$  lies in some local coordinate system  $U(x^1, \dots, x^n)$  of  $M^n$ . At  $x$  we have a basis  $\frac{\partial}{\partial x^1}, \dots, \frac{\partial}{\partial x^n}$  for  $M_x^n$ . We can express  $v = v^i \frac{\partial}{\partial x^i}$ . Then  $(x, v)$  is completely described by the  $2n$ -tuple of real numbers  $(x^1, \dots, x^n, v^1(x), \dots, v^n(x))$ .

If  $x$  also lies in the coordinate patch  $\bar{U}(\bar{x}^1, \dots, \bar{x}^n)$  the same  $(x, v) \in TM^n$  would be described by the new  $2n$ -tuple  $(\bar{x}^1, \dots, \bar{x}^n, \bar{v}^1(x), \dots, \bar{v}^n(x))$  where

$$\begin{cases} \bar{x}^i = \bar{x}^i(x^1, \dots, x^n) \\ \bar{v}^i = \frac{\partial \bar{x}^i}{\partial x^j}(x) v^j(x) \end{cases}$$

It is obvious that  $TM^n$  is a  $2n$ -dimensional differentiable manifold.

Cotangent bundle to  $M^n$  is the space  $T^*M^n$  of all covectors on  $M^n$ . A point in  $T^*M^n$  is a pair  $(x, \alpha)$ . If  $x$  is in the coordinate patch  $U(x^1, \dots, x^n)$  of  $M^n$  we have the basis  $dx^1, \dots, dx^n$  for  $(M_x^n)^*$  and  $\alpha$  can be expressed as  $\alpha = a_i dx^i$ . Then  $(x, \alpha) \in T^*M^n$  is completely described by the  $2n$ -tuple  $(x^1, \dots, x^n, a_1(x), \dots, a_n(x))$ . If  $x$  also lies in the coordinate patch  $\bar{U}(\bar{x}^1, \dots, \bar{x}^n)$  the same  $(x, \alpha) \in T^*M^n$

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would be described by the new 2n-tuple  $(\bar{x}^1, \dots, \bar{x}^n, \bar{a}_1(x), \dots, \bar{a}_n(x))$   
where

$$\begin{cases} \bar{x}^i = \bar{x}^i(x^1, \dots, x^n) \\ \bar{a}_i = \frac{dx^j}{d\bar{x}^i}(x) a^j(x) \end{cases}$$

Again it is obvious that  $T^*M^n$  is a 2n-dimensional differentiable manifold.

• Riemannian Manifold - First let's develop the idea of inner product which will allow us to associate a covector with each vector. Let  $E$  be an n-dimensional real vector space with an inner product  $\langle, \rangle$ . For each pair of vectors  $v, w \in E$  [ $\langle, \rangle: E \times E \rightarrow \mathbb{R}$ ]

(i)  $\langle v, w \rangle$  is linear in each entry when the other is held fixed

(ii)  $\langle v, w \rangle = \langle w, v \rangle$

(iii)  $\langle, \rangle$  is non-singular (i.e., if  $\langle v, w \rangle = 0 \forall w \in E \Rightarrow v = 0 \in E$ )

If  $e_1, \dots, e_n$  is a basis for  $E$  then the matrix

$$g_{ij} = \langle e_i, e_j \rangle$$

is called the (matrix of the) metric tensor. If  $v = v^i e_i$  and  $w = w^i e_i$  then

$$\langle v, w \rangle = g_{ij} v^i w^j$$

Now define the linear function  $f_v: E \rightarrow \mathbb{R}$  defined by  $f_v(w) = \langle v, w \rangle$ .  $f_v$  is obviously an element of  $E^*$

Putting  $v = f_v = v_i \sigma^i$  we have

$$v_j = v(e_j) = \langle v, e_j \rangle = \langle v^i e_i, e_j \rangle = v^i \langle e_i, e_j \rangle = v^i g_{ij}$$

Thus if  $E$  has an inner product, to each vector  $v \in E$  with components  $v^i$  we can associate a covector  $v \in E^*$  with components  $v_j = v^i g_{ij}$  having property

$$v(w) = \langle v, w \rangle$$

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Conversely to each  $v \in E^*$  with components  $v_i$  we can associate a vector  $v \in E$  with components  $v^i = v_j g^{ij}$  with  $((g^{ij})) = \text{Inverse matrix to } ((g_{ij}))$

A Riemannian metric on a manifold  $M^n$  assigns in a differentiable fashion a positive definite inner product  $\langle, \rangle$  to each tangent space  $M^n_p$ . A manifold with a Riemannian metric is called a Riemannian manifold.

In terms of a coordinate basis  $\frac{\partial}{\partial x^1}, \dots, \frac{\partial}{\partial x^n}$  we have

$$g_{ij} = \left\langle \frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j} \right\rangle$$

In an overlap with new coordinates  $\bar{x}$

$$\bar{g}_{ij} = \left\langle \frac{\partial}{\partial \bar{x}^i}, \frac{\partial}{\partial \bar{x}^j} \right\rangle = \frac{\partial x^l}{\partial \bar{x}^i} \frac{\partial x^m}{\partial \bar{x}^j} g_{lm}$$

Gradient vector - If  $M^n$  is a Riemannian metric and  $f$  is a differentiable function  $f: M^n \rightarrow \mathbb{R}$  the gradient vector  $\nabla f$  is the contravariant vector associated to the covector  $df$

$$df(w) = \langle \nabla f, w \rangle$$

In coordinates

$$(\nabla f)^i = g^{ij} \frac{\partial f}{\partial x^j}$$

- Tensors We have defined covectors from linear functionals on vectors. We may also think of vectors as linear functionals on. Although tensors will be defined over a general finite dimensional vector space  $E$ , we will usually be concerned with the case  $E = M^n_p$ . Instead of using  $e_1, \dots, e_n$  for a basis for  $E$  and  $\sigma^1, \dots, \sigma^n$  for a basis for  $E^*$  we will use  $e_i = \frac{\partial}{\partial x^i}$  and  $\sigma^i = dx^i$  with the case of primary concern in mind

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A covariant tensor of rank  $r$  is a multilinear real valued function  $Q: \underbrace{E \times \dots \times E}_{r \text{ times}} \rightarrow \mathbb{R}$ . Thus  $Q$  is a function

of  $r$ -tuples of vectors  $Q(v_{(1)}, \dots, v_{(r)})$  that is linear in each  $v_{(i)}$  when remaining entries are kept fixed. In components we have from multilinearity

$$Q(v_{(1)}, \dots, v_{(r)}) = v_{(1)}^{i_1} v_{(2)}^{i_2} \dots v_{(r)}^{i_r} Q_{i_1 i_2 \dots i_r} \quad \text{where}$$

$$Q_{i_1 i_2 \dots i_r} \equiv Q\left(\frac{\partial}{\partial x^{i_1}}, \frac{\partial}{\partial x^{i_2}}, \dots, \frac{\partial}{\partial x^{i_r}}\right)$$

$Q$  can now be written in components

$$Q = Q_{i_1 i_2 \dots i_r} dx^{i_1} \otimes dx^{i_2} \otimes \dots \otimes dx^{i_r}$$

where tensor product  $\otimes$  is defined as

$$d_{(1)} \otimes d_{(2)} \otimes \dots \otimes d_{(r)}(v_{(1)}, v_{(2)}, \dots, v_{(r)}) = d_{(1)}(v_{(1)}) d_{(2)}(v_{(2)}) \dots d_{(r)}(v_{(r)})$$

The space of all covariant tensors of rank  $r$  forms a vector space which we will write as

$$Q \in \underbrace{E^* \otimes E^* \otimes \dots \otimes E^*}_{r \text{-times}} = \otimes^r E^*$$

It has dimension  $n^r$ .

A contravariant tensor of rank  $s$  is a  $s$ -linear function on covectors  $T: \underbrace{E^* \times E^* \times \dots \times E^*}_{s \text{-times}} \rightarrow \mathbb{R}$ . In components

$$T(d_{(1)}, \dots, d_{(s)}) = d_{(1)}^{i_1} d_{(2)}^{i_2} \dots d_{(s)}^{i_s} T^{i_1 i_2 \dots i_s} \quad \text{where}$$

$$T^{i_1 i_2 \dots i_s} \equiv T(dx^{i_1}, dx^{i_2}, \dots, dx^{i_s})$$

$T$  can now be written in components

$$T = T^{i_1 i_2 \dots i_s} \frac{\partial}{\partial x^{i_1}} \otimes \frac{\partial}{\partial x^{i_2}} \otimes \dots \otimes \frac{\partial}{\partial x^{i_s}}$$



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the tensor product of vectors is defined

$$v_{(1)} \otimes v_{(2)} \otimes \dots \otimes v_{(s)}(d_{(1)}, \dots, d_{(s)}) = v_{(1)}(d_{(1)}) \dots v_{(s)}(d_{(s)})$$

The space of all contravariant tensors of rank  $s$  forms a vector space which we will write as

$$T \in \underbrace{E \otimes E \otimes \dots \otimes E}_{s\text{-times}} = \otimes^s E$$

It has dimension  $n^s$ .

A mixed tensor  $n$  times covariant,  $s$  times contravariant is a multilinear function  $W: \underbrace{E^* \times \dots \times E^*}_s \times \underbrace{E \times \dots \times E}_r \rightarrow \mathbb{R}$

In components

$$W(d_{(1)}, \dots, d_{(s)}, v_{(1)}, \dots, v_{(r)}) = d_{(1)j_1} \dots d_{(s)j_s} W^{j_1 \dots j_s}_{i_1 \dots i_r} v_{(1)}^{i_1} \dots v_{(r)}^{i_r}$$

where  $W^{j_1 \dots j_s}_{i_1 \dots i_r} \equiv W(dx^{j_1}, \dots, dx^{j_s}, \frac{\partial}{\partial x^{i_1}}, \dots, \frac{\partial}{\partial x^{i_r}})$

$W$  can now be written in components

$$W = W^{j_1 \dots j_s}_{i_1 \dots i_r} \frac{\partial}{\partial x^{j_1}} \otimes \dots \otimes \frac{\partial}{\partial x^{j_s}} \otimes dx^{i_1} \otimes \dots \otimes dx^{i_r}$$

Obviously  $W \in \otimes^r E^* \otimes^s E$  with dimension  $n^{r+s}$

Transformation Properties Under a change of basis

$$\frac{\partial}{\partial \bar{x}^i} = \frac{\partial x^k}{\partial \bar{x}^i} \frac{\partial}{\partial x^k} \quad \text{and} \quad d\bar{x}^j = \frac{\partial \bar{x}^j}{\partial x^l} dx^l \quad \text{we get by}$$

multilinearity

$$\bar{W}^{j_1 \dots j_s}_{i_1 \dots i_r} = \frac{\partial \bar{x}^{j_1}}{\partial x^{l_1}} \dots \frac{\partial \bar{x}^{j_s}}{\partial x^{l_s}} W^{l_1 \dots l_s}_{k_1 \dots k_r} \frac{\partial x^{k_1}}{\partial \bar{x}^{i_1}} \dots \frac{\partial x^{k_r}}{\partial \bar{x}^{i_r}}$$

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General Tensor Product If  $\alpha^p \in \otimes^p E^*$  and  $\beta^q \in \otimes^q E^*$  define their tensor product by

$$\alpha^p \otimes \beta^q (v_1, \dots, v_{p+q}) \equiv \alpha^p(v_1, \dots, v_p) \beta^q(v_{p+1}, \dots, v_{p+q})$$

Grassmann or Exterior Algebra

$$\text{Let } \Lambda^p E^* \equiv \underbrace{E^* \wedge E^* \wedge \dots \wedge E^*}_{p\text{-times}} \subset \otimes^p E^*$$

be the vector space of skew symmetric (alternating or antisymmetric) covariant  $p$ -tensors. Tensors such that

$$\alpha^p(v_1, \dots, v_r, \dots, v_s, \dots, v_r, \dots, v_s, \dots, v_p) = -\alpha^p(v_1, \dots, v_s, \dots, v_r, \dots, v_r, \dots, v_s, \dots, v_p)$$

The elements of  $\Lambda^p E^*$  are called exterior  $p$ -forms or simply  $p$ -forms.

Note that, if  $\alpha^p \in \Lambda^p E^*$  and  $\beta^q \in \Lambda^q E^*$  then  $\alpha^p \otimes \beta^q \in \otimes^{p+q} E^*$  but  $\alpha^p \otimes \beta^q$  will not necessarily be an element of  $\Lambda^{p+q} E^*$ . We need to define a product that is closed with respect to all forms.

It is easy to do this for  $p=q=1$

$$d^1 \wedge d^1 \equiv d^1 \otimes d^1 - d^1 \otimes d^1$$

For general  $d$  and  $\beta$  we define the exterior or Grassmann product

$$\wedge: \Lambda^p E^* \times \Lambda^q E^* \rightarrow \Lambda^{p+q} E^*$$

Let  $\alpha^p \in \Lambda^p E^*$  and  $\beta^q \in \Lambda^q E^*$  then define  $\alpha^p \wedge \beta^q \in \Lambda^{p+q} E^*$

by

$$\alpha^p \wedge \beta^q (v_{i_1}, \dots, v_{i_{p+q}})$$

$$= \sum_{\substack{k_1, \dots, k_p \\ i_1, \dots, i_p}} \sum_{\substack{l_1, \dots, l_q \\ i_{p+1}, \dots, i_{p+q}}} \delta_{i_1, \dots, i_p}^{k_1, \dots, k_p} \delta_{i_{p+1}, \dots, i_{p+q}}^{l_1, \dots, l_q} \alpha(v_{k_1}, \dots, v_{k_p}) \beta(v_{l_1}, \dots, v_{l_q})$$

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in shorthand notation

$$\alpha^p \wedge \beta^q (v_I) = \sum_{\substack{K \subset I \\ L \subset I}} \delta_I^{KL} \alpha(v_K) \beta(v_L)$$

where

$$I = (i_1, \dots, i_{p+q})$$

$$K = (k_1, \dots, k_p)$$

$$L = (l_1, \dots, l_q)$$

$$K \rightarrow k_1 < \dots < k_p$$

$$L \rightarrow l_1 < \dots < l_q$$

and

$$\delta_{\substack{i_1, \dots, i_r \\ j_1, \dots, j_r}} = \begin{cases} 1 & \text{if } (j_1, \dots, j_r) \text{ is even permutation of } (i_1, \dots, i_r) \\ -1 & \text{if } (j_1, \dots, j_r) \text{ is odd permutation of } (i_1, \dots, i_r) \\ 0 & \text{if } (j_1, \dots, j_r) \text{ is not a permutation of } (i_1, \dots, i_r) \end{cases}$$

One can now define Grassman or Exterior algebra on the space of all forms over  $E^*$

$$\Lambda^* E \equiv (\Lambda^0 E^* = \mathbb{R}) \oplus (\Lambda^1 E^* = E^*) \oplus (\Lambda^2 E^*) \oplus \dots \oplus (\Lambda^n E^*)$$

Obviously the  $\dim \Lambda^p E^* = \binom{n}{p}$  where  $\dim(E) = n$ .

It follows that

$$\dim(\Lambda^* E) = \binom{n}{0} + \binom{n}{1} + \binom{n}{2} + \dots + \binom{n}{n} = 2^n$$

It can be easily verified

$$\alpha^p \wedge \beta^q = (-1)^{pq} \beta^q \wedge \alpha^p \quad \text{"anti-commutative"}$$

$$(\alpha \wedge \beta) \wedge \gamma = \alpha \wedge (\beta \wedge \gamma) \quad \text{associative}$$

$$\alpha \wedge \beta \in \Lambda^* E \quad \text{closed}$$

Any  $p$ -form can be expanded in terms of the  $\alpha$  basis defined as follows: Let  $\sigma^1, \dots, \sigma^n$  be the dual basis to  $e_1, \dots, e_n$ .  $\alpha^p \in \Lambda^p E^*$  can now be expanded as

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$$\alpha^p = \sum_{\substack{I \\ \rightarrow}} \alpha_I \sigma^I = \sum_{i_1 < \dots < i_p} \alpha_{i_1 \dots i_p} \sigma^{i_1} \wedge \dots \wedge \sigma^{i_p}$$

where  $\alpha_I \equiv d(e_I)$

$$\sigma^I(e_J) = \delta^I_J$$

## Exterior Differentiation

Thm: There is a unique operator (exterior differentiation)

$$d: \Lambda^p M^* \rightarrow \Lambda^{p+1} M^*$$

satisfying

(i)  $d$  is additive  $d(\alpha^p + \beta^p) = d\alpha^p + d\beta^p$

(ii)  $d\alpha^0 = d$  differential of the function  $\alpha^0$

(iii)  $d(\alpha^p \wedge \beta^q) = d\alpha^p \wedge \beta^q + (-1)^p \alpha^p \wedge d\beta^q$

(iv)  $d^2 \alpha^p = d(d\alpha^p) \equiv 0 \quad (\forall \alpha)$

If  $f$  is a 0-form (function)

$$df = \frac{df}{dx^i} dx^i \quad (\text{from } df(v) = v(f))$$

If  $a_I(x)$  are functions and  $\alpha = \sum_{\substack{I \\ \rightarrow}} a_I \sigma^I$

$$d\alpha = \sum_{\substack{I \\ \rightarrow}} da_I \wedge \sigma^I$$

It can easily be shown that the pull back of a form under a map  $F: M^n \rightarrow V^r$  is again a form and this pull back commutes with both exterior differentiation and the wedge product.

(i) If  $\alpha^p \in \Lambda^p V^r$   $F^* \alpha \in \Lambda^p M^n$

(ii)  $F^*(\alpha \wedge \beta) = (F^* \alpha) \wedge (F^* \beta)$

(iii)  $F^*(d\alpha) = d(F^* \alpha)$

(13)

• Interior Product - If  $\bar{X}$  is a vector and  $\alpha^p$  is a  $p$ -form we define the interior product or contraction as:

$$i_{\bar{X}}: \Lambda^p M^* \rightarrow \Lambda^{p-1} M^*$$

$$\begin{cases} i_{\bar{X}} \alpha^0 = 0 \\ i_{\bar{X}} \alpha^1 = \alpha^1(\bar{X}) \\ i_{\bar{X}} \alpha^p(\bar{X}_2, \bar{X}_3, \dots, \bar{X}_p) = \alpha^p(\bar{X}, \bar{X}_2, \dots, \bar{X}_p) \end{cases}$$

In a basis, if  $\alpha = a_I \sigma^I$ ,

$$i_{\bar{X}} \alpha = b_J \sigma^J = \bar{X}^i a_{iJ} \sigma^{iJ}$$

Also  $i_{\bar{X}}$  can be shown to be an anti-derivation that is

$$i_{\bar{X}}(\alpha^p \wedge \beta^q) = (i_{\bar{X}} \alpha^p) \wedge \beta^q + (-)^p \alpha^p \wedge (i_{\bar{X}} \beta^q)$$

## Lie Derivative

We must first define the flow  $\phi_t$  generated by a vector field  $\bar{X}$  on a manifold  $M$ .

A vector field  $\bar{X}$  on  $U \subset M^n$  assigns in a differentiable manner a vector  $v_p$  to each  $p \in U$ . In coordinates any vector field on a patch  $(U, x)$  has a local expression

$$\bar{X} = v^i(x) \frac{\partial}{\partial x^i}$$

A flow is a 1-parameter family of maps

$$\phi_t: M \rightarrow M$$

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such that (i)  $\phi_s(\phi_t(p)) = \phi_{s+t}(p) = \phi_t(\phi_s(p))$   $\left( \begin{array}{l} \forall s, t \in \mathbb{R} \\ \forall p \in M \end{array} \right)$   
 (ii)  $\phi_{-t}(\phi_t(p)) = p$   
 (iii)  $\phi_t$  is a diffeomorphism

associated with  $\phi_t$  is a time independent velocity vector field

$$V_p = \left[ \frac{d}{dt} (\phi_t p) \right]_{t=0}$$

$$\text{or } V_p(f) = \left[ \frac{d}{dt} f(\phi_t p) \right]_{t=0} \text{ for } f: M \rightarrow \mathbb{R}$$

Can be shown that for each differential vector field  $X$  one can associate a flow  $\phi_t$  such that  $X$  is the velocity field of  $\phi_t$ .

Given two vector fields  $X$  and  $Y$  on  $M$  the Lie bracket or Lie derivative of  $Y$  with respect to  $X$  is the mapping

$$\mathcal{L}_X : X_0^* M \rightarrow X_0^* M \quad \text{where } X_0^* M = \prod_{x \in M} M_x$$

such that vector field is differentiable

$$\begin{aligned} (\mathcal{L}_X Y)_x &= \lim_{t \rightarrow 0} \frac{Y_{\phi_t x} - (\phi_t)_* Y_x}{t} \quad \text{where } \phi_t \text{ is flow generated by } X \\ & \quad x \in M \quad X, Y \in X_0^* M \\ &= \lim_{t \rightarrow 0} \frac{(\phi_{-t})_* Y_{\phi_t x} - Y_x}{t} = \frac{d}{dt} \left[ (\phi_{-t})_* Y_{\phi_t x} \right]_{t=0} \end{aligned}$$

Can be shown that

$$\mathcal{L}_X Y(f) = X(Y(f)) - Y(X(f)) = [X, Y](f)$$

(15)

In local coordinates where  $X = X^i \frac{\partial}{\partial x^i}$   $\Upsilon = \Upsilon^i \frac{\partial}{\partial x^i}$

$$(\mathcal{L}_X \Upsilon)^i = [\Upsilon, X]^i = X^j \frac{\partial \Upsilon^i}{\partial x^j} - \Upsilon^j \frac{\partial X^i}{\partial x^j}$$

Lie derivative of forms If  $X$  is a vector field (with local flow  $\phi_t$ ) and  $f$  is a function  $f: M \rightarrow \mathbb{R}$  we will define the Lie derivative of  $f$  with respect to  $X$  at  $p$  by  $\mathcal{L}_X: X_0^* M^* \rightarrow X_0^* M^*$

$$\mathcal{L}_X(f) \equiv X_p(f) = \left. \frac{d}{dt} f(\phi_t p) \right|_{t=0} = \left. \frac{d}{dt} (\phi_t^* f) \right|_{t=0}$$

If  $\alpha^p$  is  $p$ -form we define

$$\mathcal{L}_X \alpha^p \equiv \left. \frac{d}{dt} (\phi_t^* \alpha^p) \right|_{t=0} = \lim_{t \rightarrow 0} \frac{\phi_t^* \alpha - \alpha}{t}$$

In particular if  $v_1, \dots, v_p$  are invariant under the flow  $\phi_t$  generated by  $X$  (i.e.,  $\mathcal{L}_X v_i = 0$ ) then

$$\mathcal{L}_X(\alpha^p)(v_1, \dots, v_p) = \left. \frac{d}{dt} [\alpha_{\phi_t x}^p(v_1, \dots, v_p)] \right|_{t=0}$$

Properties of  $\mathcal{L}_X$  that are easily verified

- (i)  $\mathcal{L}_X(\alpha^p \wedge \beta^q) = (\mathcal{L}_X \alpha^p) \wedge \beta^q + \alpha^p \wedge (\mathcal{L}_X \beta^q)$  "derivation"
- (ii)  $\mathcal{L}_X \circ d = d \circ \mathcal{L}_X$
- (iii)  $\mathcal{L}_X = i_X \circ d + d \circ i_X$  "Cartan's Formula"
- (iv)  $X(\alpha^p(u_1, \dots, u_p)) = (\mathcal{L}_X \alpha^p)(u_1, \dots, u_p) + \sum_r \alpha^p(u_1, \dots, [X, u_r], \dots, u_p)$   
 $X, u_1, \dots, u_p$  vector field "Leibnitz rule"
- (v)  $d\alpha^p(u_0, \dots, u_p) = \sum_{r=0}^p (-1)^r i_{u_r} \alpha(u_0, \dots, \hat{u}_r, \dots, u_p)$   
 $+ \sum_{r < s} (-1)^{r+s} \alpha^p([u_r, u_s], u_0, \dots, \hat{u}_r, \dots, \hat{u}_s, \dots, u_p)$

(note: vectors on left side vector fields on right.  $\wedge$  means outer product)

## Orientation and Pseudo Forms

Let  $e = (e_1, \dots, e_n)$  and  $f = (f_1, \dots, f_n)$  be two bases of a vector space  $E$ . We can then uniquely write  $f = eA$  (i.e.,  $f_i = e_j A^j_i$ ) for some non-singular matrix  $A$ . If  $\det A > 0$  (resp.  $< 0$ ) we say that  $e$  and  $f$  have the same (resp. opposite) orientation.

We shall say that the entire manifold  $M^n$  is orientable provided we can cover  $M^n$  by coordinate patches  $\{U_\alpha\}$  such that in each overlap  $U_\alpha \cap U_\beta$  the Jacobian determinant

$$\frac{d(x^1_\beta, \dots, x^n_\beta)}{d(x^1_\alpha, \dots, x^n_\alpha)} > 0$$

Obviously this says the two orientations of  $M^n$  coming from  $U_\alpha$  and  $U_\beta$  agree since

$$e_i = \frac{\partial}{\partial x^i_\alpha} = \sum_j \frac{\partial x^j_\beta}{\partial x^i_\alpha} \frac{\partial}{\partial x^j_\beta} = \sum_j A^j_i f_j$$

The forms previously discussed were independent of the notion of orientability;  $d^p(v_1, \dots, v_p)$  is a number which has nothing to do with the orientation of the particular tangent space in which  $v_1, \dots, v_n$  live.

There is a different notion of exterior form which is dependant on the orientation



(17)

A pseudo form of degree  $p$  on a manifold  $M^n$  assigns to each orientation,  $\sigma$ , of the tangent space  $M_x^n$  a exterior  $p$ -form in such a way that

$$d^p_{-\sigma} = -d^p_{\sigma}$$

Note: an example of a commonly encountered pseudo-form is the volume form

$$\text{vol}_{\sigma}^3 = \alpha(x) \sqrt{\det g_{ij}(x)} dx^1 \wedge dx^2 \wedge dx^3 \quad (\text{Riemannian Manifold})$$

$n=3$

where

$$\alpha(x) = \begin{cases} +1 & \text{if } \frac{\partial}{\partial x^1} \frac{\partial}{\partial x^2} \frac{\partial}{\partial x^3} \text{ is positively oriented w.r.t } \sigma \\ -1 & \text{if } \frac{\partial}{\partial x^1} \frac{\partial}{\partial x^2} \frac{\partial}{\partial x^3} \text{ is negatively oriented w.r.t } \sigma \end{cases}$$

The reason for all this nit picking about orientation and pseudo forms will become apparent with the development of integration of forms over manifolds.

## Integration of Forms over Oriented Manifolds

Partition of Unity - Given a finite coordinate covering  $\{U_{\alpha}, \varphi_{\alpha}\}$  ( $\alpha=1, \dots, N$ ) a partition of unity will exhibit  $N$  real valued functions  $f_{\alpha} : V^p \rightarrow \mathbb{R}$  that are differentiable and have the following properties

- (i)  $f_{\alpha}(x) \geq 0$  ( $\alpha=1, \dots, N$ )
- (ii) The support of  $f_{\alpha}$  is a closed subset  $\subset U_{\alpha}$   
in particular  $f_{\alpha}(p) = 0$  for  $p \notin U_{\alpha}$
- (iii)  $\sum_{\alpha} f_{\alpha}(p) = 1 \quad \forall p \in V^p$

the support of a function is the closure of a set of points at which the function is not zero

It can be shown that a partition of unity can always be constructed for every coordinate cover.

Integration of forms - Let  $M^n$  be a manifold.

Let  $V^p$  be an oriented submanifold (perhaps with boundary) of  $M^n$ . We can think of  $V^p$  as a manifold in its own right. We then have the inclusion map

$$i: V^p \rightarrow M^n$$

where  $i(x)$  says think of the point  $x \in V^p$  as a point in  $M^n$ . If  $u^1, \dots, u^p$  are local coordinates for  $V^p$  and  $x^1, \dots, x^n$  those for  $M^n$  then  $i$  is simply given by  $x^i = x^i(u^1, \dots, u^p)$

Let  $\omega^p$  be a  $p$ -form on  $M^n$ . We may restrict the form to  $V^p$  by taking the pull back  $i^*\omega^p$ . Now let  $\{U_\alpha, \varphi_\alpha\}$  be a finite covering of  $V^p$  by coordinate charts that are positively oriented for the orientation of  $V^p$  given. We then define

$$\begin{aligned} \int_{V^p} \omega^p &\equiv \int_{V^p} i^*\omega^p = \int_{V^p} (i^*\omega^p) \sum_\alpha f_\alpha = \sum_\alpha \int_{V^p} f_\alpha i^*\omega^p \\ &= \sum_\alpha \int_{U_\alpha} f_\alpha i^*\omega^p \equiv \sum_\alpha \int_{U_\alpha} \omega_\alpha^p \quad \text{where } \omega_\alpha^p = f_\alpha i^*\omega^p \\ &= \sum_\alpha \int_{\varphi_\alpha(U_\alpha)_0} \varphi_\alpha^* \omega_\alpha^p = \sum_\alpha \int_{\varphi_\alpha(U_\alpha)_0 \subset \mathbb{R}^p} f_\alpha(\varphi_\alpha^{-1}(u)) (i \circ \varphi_\alpha^{-1})^* \omega^p \\ &= \sum_{\varphi_\alpha(U_\alpha)_0} \int f_\alpha(u) (i \circ \varphi_\alpha^{-1})^* \omega^p \left( \frac{\partial}{\partial u^1}, \dots, \frac{\partial}{\partial u^p} \right) du^1 \dots du^p \\ &\equiv \sum_\alpha \int_{|\varphi_\alpha(U_\alpha)|} \sigma(u) f_\alpha(u) \omega^p \left( (i \circ \varphi_\alpha^{-1})_* \frac{\partial}{\partial u^1}, \dots, (i \circ \varphi_\alpha^{-1})_* \frac{\partial}{\partial u^p} \right) du^1 \dots du^p \end{aligned}$$

(note: this last integral is the normal multiple Riemann integral)

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where  $i \circ \varphi_\alpha^{-1} : \mathbb{R}^p \rightarrow M^n$   
 $f_\alpha(p) : U_\alpha \subset V^p \rightarrow \mathbb{R}$   
 $f_\alpha(u) = f_\alpha(\varphi_\alpha^{-1}(u)) : \mathbb{R}^p \rightarrow \mathbb{R}$

$|\varphi_\alpha(U_\alpha)|$  will be  $\varphi_\alpha(U_\alpha)$  without orientation  
orientation for  $\varphi_\alpha(U_\alpha)$  will be induced by the  
orientation for  $U_\alpha$  which is induced by the  
orientation for  $V^p$

Notice that this integration does not depend upon the  
orientation of  $M^n$ , only on the orientation of  $V^p$ . If the  
orientation of  $V^p$  is reversed so is the value of the  
integral.

### Integration of pseudo-forms

For pseudo-forms either an orientation must be given  
for  $M^n$  (i.e.,  $M^n$  must be orientable) or some orientation  
must be given for each coordinate patch. Any orientation  
can be given to each coordinate patch because of the  
fact that the value of the integral of a pseudo-form  
is independent of the orientation of  $M^n$  or  $V^p$

$$\int_{V_0^p} \omega_0^p = - \int_{V_0^p} -\omega_0^p = \int_{V_0^p} \omega_0^p$$

This adds some technical details with regard to defining  
an orientation for each map but leads to the advantage  
that pseudo forms can be integrated over  $V^p$  that  
are not orientable where true forms must be  
integrated over orientable  $V^p$  since the orientation  
matters for them.

## ★ Stokes's Thm ★

A manifold with boundary is the same as a manifold except that charts  $(U, \varphi)$  of the following form are allowed

$$\varphi: U \rightarrow \{x \in \mathbb{R}^n : x^n \geq 0\}$$

the boundary of  $M$ ,  $\partial M$ , is the set of points

$$\partial M = \{x \in M : \varphi_\alpha(x) = 0 \text{ for some } \alpha\}$$

Now for one of the main results of exterior calculus.

Thm (Stokes's): Let  $V \subset M^n$  be a compact submanifold-with-boundary. Let  $\omega^{p-1}$  be a continuously differentiable true ( $V$  must be orientable) or pseudo  $p-1$  form on  $M^n$ . Then

$$\int_V d\omega^{p-1} = \int_{\partial V} \omega^{p-1}$$

## Differentiation of Integrals

Time-independent Let  $d^p$  be a  $p$ -form (time-independent),  $C_p$  a  $p$  dimensional submanifold (perhaps with boundary  $\partial C_p$ ) of a manifold  $M^n$ . We shall consider a "variation" of  $C_p$  arising as follows. We suppose that there is a flow  $\phi_t: M^n \rightarrow M^n$ .  $C_p(t) = \phi_t C_p$  We are interested in the variation of the integral  $I(t) \equiv \int_{C_p(t)} d^p = \int_{\phi_t C_p} d^p = \int_{C_p} \phi_t^* d^p$

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Differentiating,

$$\begin{aligned} I'(t) &= \lim_{h \rightarrow 0} \frac{1}{h} [I(t+h) - I(t)] \\ &= \lim_{h \rightarrow 0} \frac{1}{h} \left[ \int_{C_P} \phi_{t+h}^* \alpha^p - \int_{C_P} \phi_t^* \alpha^p \right] = \lim_{h \rightarrow 0} \frac{1}{h} \int_{C_P} \phi_t^* (\phi_h^* \alpha^p - \alpha^p) \\ &= \lim_{h \rightarrow 0} \frac{1}{h} \int_{\phi_t C_P} \phi_h^* \alpha^p - \alpha^p = \int_{\phi_t C_P} \lim_{h \rightarrow 0} \frac{1}{h} (\phi_h^* \alpha^p - \alpha^p) \end{aligned}$$

$$\Rightarrow \boxed{\frac{d}{dt} \int_{C_P(t)} \alpha^p = \int_{C_P(t)} \mathcal{L}_{\vec{v}} \alpha^p \stackrel{\text{Cartan's formula}}{=} \int_{C_P(t)} i_{\vec{v}} d\alpha^p + \int_{dC_P(t)} i_{\vec{v}} \alpha^p}$$

Time-dependant Let  $\alpha^p$  be a time dependant  $p$ -form on  $M^n$   $\alpha^p = a_I(x, t) \sigma^I$ . Suppose we have a time dependant "flow" of fluid in  $M^n$  with velocity flow  $\vec{v}(t, x) = v^i(t, x) \frac{\partial}{\partial x^i}$ . We wish to calculate

$$\frac{d}{dt} \int_{C_P(t)} \alpha^p$$

We pass to  $\mathbb{R} \times M^n$ .  $\alpha^p = \sum_{i_1 < \dots < i_p} a_{i_1 \dots i_p} dx^{i_1} \wedge \dots \wedge dx^{i_p}$  (no  $dt$  terms)

and  $X = \frac{\partial}{\partial t} + v^i \frac{\partial}{\partial x^i}$  generating the flow  $\psi_t$ .

Let  $d = dt \wedge \frac{\partial}{\partial t} + \vec{d}$  with  $\vec{d} = dx^i \wedge \frac{\partial}{\partial x^i}$ . Using the formula for the time independant case yields the formula

$$\begin{aligned} \frac{d}{dt} \int_{C_P(t)} \alpha^p &= \int_{C_P(t)} \frac{d\alpha^p}{dt} + i_{\vec{v}} \vec{d} \alpha^p + \vec{d} i_{\vec{v}} \alpha^p \\ &= \int_{C_P(t)} \frac{d\alpha^p}{dt} + i_{\vec{v}} \vec{d} \alpha^p + \int_{dC_P(t)} i_{\vec{v}} \alpha^p \quad (\text{space only}) \end{aligned}$$

# Vector Algebra and Calculus

First lets associate a true vector  $\vec{A}$  with a true 1-form,  $\alpha^1$ , and a pseudo 2-form,  $\beta_0^2$ , the following way:

$$\alpha^1 = \langle \cdot, \vec{A} \rangle$$

$$\beta_0^2 = i_{\vec{A}} \text{vol}_0^3 = \text{ocw} \sqrt{g_{\text{cw}}} (a^1 du^1 \wedge du^3 + a^2 du^3 \wedge du^1 + a^3 du^1 \wedge du^2)$$

for a Riemannian manifold and vector  $\vec{A} = a^i \frac{\partial}{\partial x^i}$ .

Conversly, one can associate a pseudo-vector  $\vec{B}_0$  with a pseudo 1-form and a true 2 form, the following way:

$$\beta_0^1 = \langle \cdot, \vec{B}_0 \rangle$$

$$\beta_0^2 = i_{\vec{B}_0} \text{vol}_0^3$$

Given a pair of true vectors  $\vec{A}$  and  $\vec{B}$  one can express there dot product as a 0-form or a 3-form in  $\mathbb{R}^3$

$$i_{\vec{B}} \alpha^1 = \vec{A} \cdot \vec{B} \quad \text{true 0-form}$$

$$\alpha^1 \wedge \beta_0^2 = (\vec{A} \cdot \vec{B}) \text{vol}_0^3 \quad \text{pseudo 3-form}$$

The pseudo vector,  $\vec{A} \times \vec{B}$  can be associated with a 2-form and 1-form in  $\mathbb{R}^3$

$$\alpha^1 \wedge \beta_0^1 = i_{\vec{A} \times \vec{B}} \text{vol}_0^3 \quad \text{true 2-form}$$

$$-i_{\vec{A}} \beta_0^2 = \langle \cdot, \vec{A} \times \vec{B} \rangle \quad \text{pseudo 1-form}$$

The vector triple product can be associated with a 0-form or 3-form in  $\mathbb{R}^3$

$$\alpha^1 \wedge \beta_0^1 \wedge \beta_0^1 = (\vec{A} \times \vec{B}) \cdot \vec{C} \text{vol}_0^3 \quad \text{true 3-form}$$

$$-i_{\vec{A}} \beta_0^2 (\vec{C}) = (\vec{A} \times \vec{B}) \cdot \vec{C} \quad \text{pseudo 0-form}$$

(23)

Vector identities can now be easily proved via forms, the only decision being whether to use 1-forms or 2-forms.

$$\begin{aligned}\langle \cdot, \vec{A} \times (\vec{B} \times \vec{C}) \rangle &= -i_{\vec{A}}(\beta^1 \wedge \gamma^1) = (-i_{\vec{A}} \beta^1) \gamma^1 + \beta^1 i_{\vec{A}} \gamma^1 \\ &= (-\vec{A} \cdot \vec{B}) \langle \cdot, \vec{C} \rangle + \langle \cdot, \vec{B} \rangle (\vec{A} \cdot \vec{C}) \\ &= \langle \cdot, \vec{B}(\vec{A} \cdot \vec{C}) - \vec{C}(\vec{A} \cdot \vec{B}) \rangle \\ \Rightarrow \vec{A} \times (\vec{B} \times \vec{C}) &= \vec{B}(\vec{A} \cdot \vec{C}) - \vec{C}(\vec{A} \cdot \vec{B})\end{aligned}$$

Now for calculus. By definition and problem #

$$df = \langle \cdot, \nabla f \rangle \quad \text{true 1-form (if time-scalar)}$$

$$d\alpha^1 = i_{\nabla \times \vec{A}} \text{vol}_0^2 \quad \text{true 2-form}$$

$$d\beta^2 = (\nabla \cdot \vec{B}) \text{vol}_0^3 \quad \text{pseudo 3-form}$$

To compute  $\text{div } \vec{B}$

$$(\nabla \cdot \vec{B}) \text{vol}_0 = d(i_{\vec{B}} \text{vol}_0) = \left\{ \frac{1}{\sqrt{g}} \sum_i \frac{\partial}{\partial u^i} (\sqrt{g} b^i) \right\} \text{vol}_0$$

$$\Rightarrow \nabla \cdot \vec{B} = \frac{1}{\sqrt{g}} \sum_i \frac{\partial}{\partial u^i} (\sqrt{g} b^i)$$

To compute a vector identity

$$[\nabla \cdot (\vec{A} \times \vec{B})] \text{vol}_0 = d(d^1 \wedge \beta^1) = d\alpha^1 \wedge \beta^1 - d^1 \wedge d\beta^1$$

$$= \underbrace{(i_{\nabla \times \vec{A}} \text{vol}_0) \wedge \beta^1}_{\text{true 2-form for } (\nabla \times \vec{A})_0} - d^1 \wedge \underbrace{i_{\nabla \times \vec{B}} \text{vol}_0}_{\text{true 2-form for } (\nabla \times \vec{B})_0}$$

$$= [\vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B})] \text{vol}_0$$

$$\Rightarrow \nabla \cdot (\vec{A} \times \vec{B}) = (\nabla \times \vec{A}) \cdot \vec{B} - (\nabla \times \vec{B}) \cdot \vec{A}$$

(24)

One last thing to note, by Cartan's formula

$$\begin{aligned} \mathcal{L}_X \text{vol}^3 &= \underbrace{i_X d \text{vol}^3}_=0 + d i_X \text{vol}^3 = d(i_X \text{vol}^3) \\ &= (\nabla \cdot X) \text{vol}^3 \end{aligned}$$

To compute  $\nabla \times (\vec{A} \times \vec{B})$  do the following

$$i_{\nabla \times (\vec{A} \times \vec{B})} \text{vol}_0^3 = d(i_{\vec{B}} d^2) \xrightarrow{\text{Cartan's formula}} \mathcal{L}_{\vec{B}} d^2 - i_{\vec{B}} d d^2$$

$$= \mathcal{L}_{\vec{B}} d^2 - i_{\vec{B}} d(i_{\vec{A}} \text{vol}^3)$$

$$= \mathcal{L}_{\vec{B}} d^2 - i_{\vec{B}} (\nabla \cdot \vec{A} \text{vol}^3)$$

$$= \mathcal{L}_{\vec{B}} d^2 - (\nabla \cdot \vec{A}) i_{\vec{B}} \text{vol}^3$$

$$= \mathcal{L}_{\vec{B}} i_{\vec{A}} \text{vol}^2 - (\nabla \cdot \vec{A}) i_{\vec{B}} \text{vol}^3$$

$$\stackrel{\text{Prob 17}}{=} i_{\vec{A}} \mathcal{L}_{\vec{B}} \text{vol}^3 + i[\vec{B}, \vec{A}] \text{vol}^3 - (\nabla \cdot \vec{A}) i_{\vec{B}} \text{vol}^3$$

$$= (\nabla \cdot \vec{B}) i_{\vec{A}} \text{vol}^3 + i[\vec{B}, \vec{A}] \text{vol}^3 - (\nabla \cdot \vec{A}) i_{\vec{B}} \text{vol}^3$$

$$= i \left( (\nabla \cdot \vec{B}) \vec{A} + [\vec{B}, \vec{A}] - (\nabla \cdot \vec{A}) \vec{B} \right) \text{vol}^2$$

$$\Rightarrow \nabla \times (\vec{A} \times \vec{B}) = (\nabla \cdot \vec{B}) \vec{A} - (\nabla \cdot \vec{A}) \vec{B} + [\vec{B}, \vec{A}]$$

to relate to conventional formulas

$$[\vec{B}, \vec{A}]^i = B^j \frac{\partial A^i}{\partial x^j} - A^j \frac{\partial B^i}{\partial x^j} = \vec{B} \cdot \nabla A^i - \vec{A} \cdot \nabla B^i$$

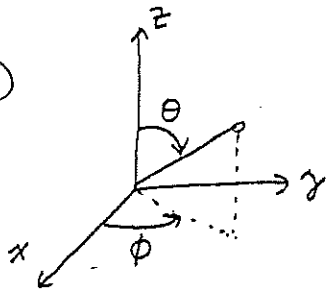
$$[\vec{B}, \vec{A}] \equiv \vec{B} \cdot \nabla \vec{A} - \vec{A} \cdot \nabla \vec{B}$$

$$\Rightarrow \nabla \times (\vec{A} \times \vec{B}) = (\nabla \cdot \vec{B}) \vec{A} - (\nabla \cdot \vec{A}) \vec{B} + \vec{B} \cdot \nabla \vec{A} - \vec{A} \cdot \nabla \vec{B}$$



# Problems on Exterior Calculus

①



In  $\mathbb{R}^3$  let  $x^1 = x$   $x^2 = y$   $x^3 = z$   
 let  $\bar{x}^1 = r$   $\bar{x}^2 = \theta$   $\bar{x}^3 = \phi$

(a) compute  $\bar{g}_{ij}$  in coordinates  $\bar{x}$

$$\bar{g}_{1\theta} \equiv \bar{g}_{\theta 1} = \left\langle \frac{d}{dr}, \frac{d}{d\theta} \right\rangle$$

(use: chain rule  $\frac{d}{dr} = \frac{\partial x}{\partial r} \frac{d}{dx} + \dots$ )

(b) Compute coefficients  $(\nabla f)^i$  in

$$\nabla f = (\nabla f)^r \frac{d}{dr} + (\nabla f)^\theta \frac{d}{d\theta} + (\nabla f)^\phi \frac{d}{d\phi}$$

(c) Verify that  $\frac{d}{dr}$   $\frac{d}{d\theta}$   $\frac{d}{d\phi}$  are orthogonal but not unit vectors

(d) Put  $\hat{e}_r = \frac{d}{dr} / \|\frac{d}{dr}\|$   $\hat{e}_\theta = \frac{d}{d\theta} / \|\frac{d}{d\theta}\|$   $\hat{e}_\phi = \frac{d}{d\phi} / \|\frac{d}{d\phi}\|$

and write  $\nabla f = (\nabla f)^r \hat{e}_r + (\nabla f)^\theta \hat{e}_\theta + (\nabla f)^\phi \hat{e}_\phi$

(? Why do we need  $\nabla f$  when  $df$  is so easy?)

② Compute in terms of a basis:

(a)  $d^1 \wedge \beta^1$

(b)  $d^1 \wedge \gamma^2$

(c)  $d^1 \wedge \beta^1 \wedge \alpha^1$

③ Compute  $(d^1 \wedge \beta^1)_I$  (ie components of  $d^1 \wedge \beta^1$  in terms of a basis)

④ Compute in terms of  $dx$   $dy$   $dz$  in  $\mathbb{R}^3$

(a)  $dd^0$

(b)  $dd^1$

(c)  $dd^2$

(d)  $d(dd^1)$

(2)

⑤ Show that

$$(a) F^*(dx \wedge dz) = \frac{\partial(x, z)}{\partial(u, v)} du \wedge dv$$

$$\text{for } F: \mathbb{R}^2 \rightarrow \mathbb{R}^2 \\ (x, z) \rightarrow (u, v)$$

$$(b) F^*(dx \wedge dy \wedge dz) = \frac{\partial(x, y, z)}{\partial(u, v, w)} du \wedge dv \wedge dw$$

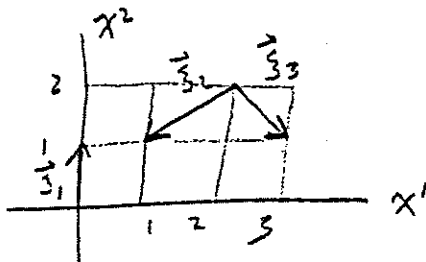
$$F: \mathbb{R}^3 \rightarrow \mathbb{R}^3 \\ (x, y, z) \rightarrow (u, v, w)$$

⑥ Calculate the value of the forms

$$\omega_1 = dx^1$$

$$\omega_2 = x^1 dx^2$$

$$\omega_3 = dr^2 \quad (r^2 = (x^1)^2 + (x^2)^2)$$

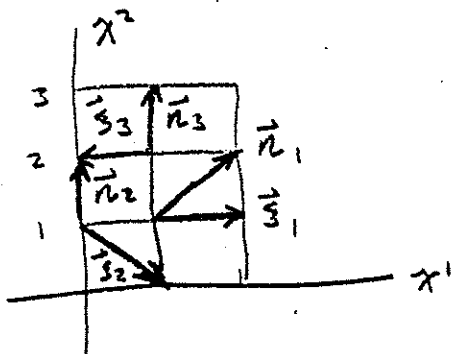
on the vectors  $\vec{\xi}_1$ ,  $\vec{\xi}_2$  and  $\vec{\xi}_3$ 

⑦ Calculate the values of the forms

$$\omega_1 = dx^1 \wedge dx^2$$

$$\omega_2 = x^1 dx^1 \wedge dx^2 - x^2 dx^2 \wedge dx^1$$

$$\omega_3 = r dr \wedge d\varphi \quad (x^1 = r \cos \varphi, x^2 = r \sin \varphi)$$

on the pairs of vectors  $(\vec{\xi}_1, \vec{\eta}_1)$ ,  $(\vec{\xi}_2, \vec{\eta}_2)$ ,  $(\vec{\xi}_3, \vec{\eta}_3)$ 

(3)

⑧ Calculate the value of the forms

$$w_1 = dx^2 \wedge dx^3$$

$$w_2 = x^1 dx^3 \wedge dx^2$$

$$w_3 = dx^3 \wedge dr^2 \quad r^2 = (x^1)^2 + (x^2)^2 + (x^3)^2$$

on  $\vec{s} = (1, 1, 1)$ ,  $\vec{n} = (1, 2, 3)$  at point  $\vec{x} = (2, 0, 0)$

⑨ Use problem #1 to write down 'vol<sup>3</sup>' for  $\mathbb{R}^3$  in spherical coordinates  $r, \theta, \varphi$

⑩ For the case  $M^n = \mathbb{R}^3$

(a) line integral

$$\varphi(\mathcal{C}) = [a, b] \subset \mathbb{R}$$

$$i: \mathcal{C} \rightarrow \mathbb{R}^3$$

$$\varphi: \mathcal{C} \rightarrow \mathbb{R}$$

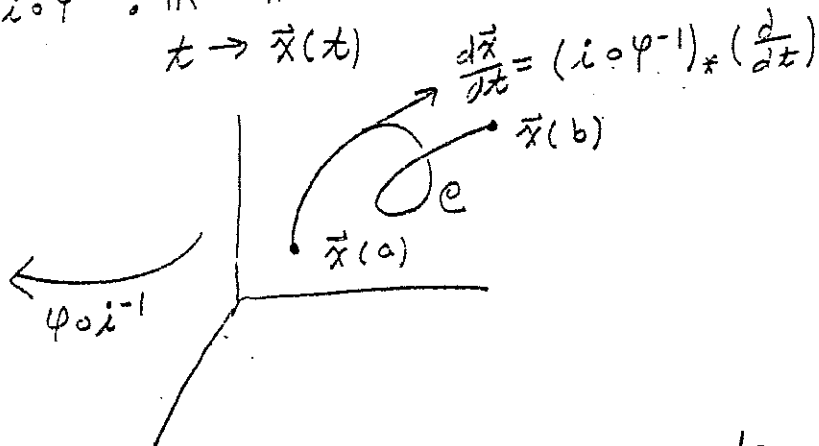
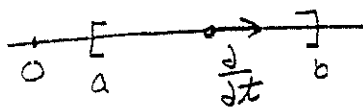
$$i \circ \varphi^{-1}: \mathbb{R} \rightarrow \mathbb{R}^3$$

$$t \rightarrow \vec{x}(t)$$

simple atlas  $\{\mathcal{C} = [a, b]\}$

$$\frac{d\vec{x}}{dt} = (i \circ \varphi^{-1})_* \left( \frac{d}{dt} \right)$$

$$\frac{d\vec{x}}{dt} = \frac{dx^i}{dt} \frac{\partial}{\partial x^i}$$



With a 1-form  $d^1 = d_i dx^i$  on  $\mathbb{R}^3$ , one can associate a contravariant vector  $\vec{a} = a^i \frac{\partial}{\partial x^i}$  since  $\mathbb{R}^3$  is Riemannian,  $d_i = g_{ij} a^j$ . Therefore given a vector field  $\vec{a}$  on  $\mathbb{R}^3$  one can construct  $d^1$ .

Show that:

$$\int_{\mathcal{C}} d^1 = \int_a^b d^1 \left( \frac{d\vec{x}}{dt} \right) dt = \int_a^b \langle \vec{a}, \frac{d\vec{x}}{dt} \rangle dt \equiv \int_{\mathcal{C}} \vec{a} \cdot d\vec{x}$$

(4)

(b) surface integral atlas  $\{(S=U, \varphi)\}$ 

$$\varphi(S) \subset \mathbb{R}^2$$

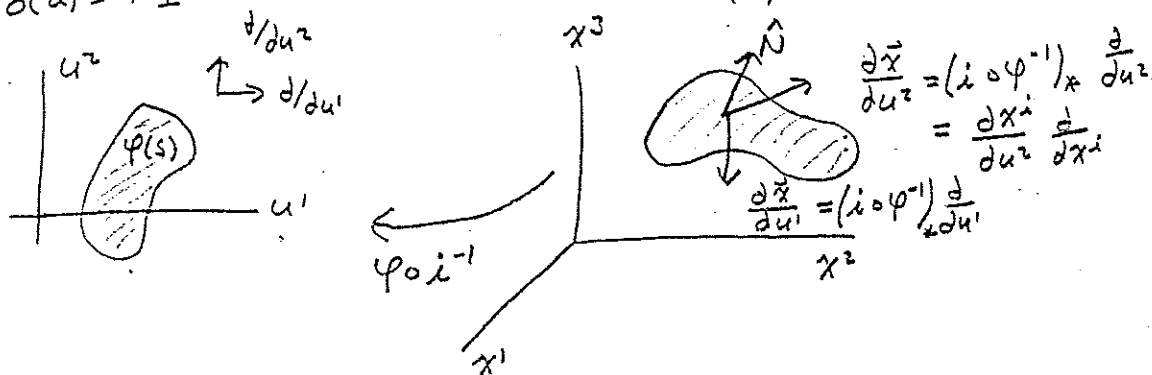
Suppose that coordinates  $u^1, u^2$  are chosen so that  $\alpha(u) = +1$

$$i: S \rightarrow \mathbb{R}^3$$

$$\varphi: S \rightarrow \mathbb{R}^2$$

$$i \circ \varphi^{-1}: \mathbb{R}^2 \rightarrow \mathbb{R}^3$$

$$(u^1, u^2) \rightarrow \vec{x}(u^1, u^2)$$



Let a 2-form  $\beta^2 = b_1 dx^2 \wedge dx^3 + b_2 dx^3 \wedge dx^1 + b_3 dx^1 \wedge dx^2$  be given on  $\mathbb{R}^3$ . One can associate a <sup>unique</sup> pseudo-vector  $\vec{B}_0$  with  $\beta^2$  by  $\beta^2 = i_{\vec{B}_0} \text{vol}_0^3$  for Riemannian  $\mathbb{R}^3$  (where  $\text{vol}_0^3 = \alpha(x) \sqrt{g(x)} dx^1 \wedge dx^2 \wedge dx^3 = \alpha(x) dx^1 \wedge dx^2 \wedge dx^3$ ). Therefore given a pseudo-vector field on  $\mathbb{R}^3$ ,  $\vec{B}_0$ , one can construct  $\beta^2$ .

Show that:

$$\int_S \beta^2 = \int_{\varphi(S)} \left[ b_1 \frac{d(x^2, x^3)}{d(u^1, u^2)} + b_2 \frac{d(x^3, x^1)}{d(u^1, u^2)} + b_3 \frac{d(x^1, x^2)}{d(u^1, u^2)} \right] du^1 du^2$$

$$= \int_{\varphi(S)} \langle \vec{B}_0, \hat{N} \rangle \text{vol}_0^3 \left( \hat{N}, \frac{d\vec{x}}{du^1}, \frac{d\vec{x}}{du^2} \right) du^1 du^2$$

$$\equiv \int \langle \vec{B}_0, \hat{N} \rangle dS \quad \text{for } \alpha \left( \hat{N}, \frac{d\vec{x}}{du^1}, \frac{d\vec{x}}{du^2} \right) = +1$$

① As a follow up to 10(b), we have seen that the "area form" for a parameterized surface in  $\mathbb{R}^3$  is the pull-back  $(i \circ \varphi^{-1})^*$  of the two form

$$i_{\hat{N}} \text{vol}^3 = dA$$

(5)

For  $i \circ \varphi^{-1}: \mathbb{R}^2 \rightarrow \mathbb{R}^3$ 

$$(u^1, u^2) \rightarrow (u^1, u^2, f(u^1, u^2))$$

One recovers the classical expression for the area

$$dS = (i \circ \varphi^{-1})_* dA \left( \frac{\partial}{\partial u^1}, \frac{\partial}{\partial u^2} \right) du^1 du^2$$

$$= \sqrt{\left( \frac{\partial f}{\partial u^1} \right)^2 + \left( \frac{\partial f}{\partial u^2} \right)^2 + 1} du^1 du^2 \quad \circ \left( \hat{N}, \frac{\partial \vec{x}}{\partial u^1}, \frac{\partial \vec{x}}{\partial u^2} \right) = +1$$

(note:  $\hat{N} = \vec{n} / \|\vec{n}\|$  where  $\vec{n} = \frac{\partial \vec{x}}{\partial u^1} \times \frac{\partial \vec{x}}{\partial u^2}$ )

(12) Apply Stokes's thm to:

(a)  $\int_{\mathbb{C} \subset \mathbb{R}^3} dd^0$  to derive Fundamental thm of Calculus(b)  $\int_{S \subset \mathbb{R}^3} dd^1$  to derive classical Stokes's Thm(c)  $\int_{V \subset \mathbb{R}^3} dd^2$  to derive divergence thm(13) Show  $\frac{d}{dt} \int_{U^n} \text{vol}^n = \int_{U^n} (\nabla \cdot \vec{v}) \text{vol}^n = \int_{\partial U^n} \langle \vec{v}, \hat{N} \rangle \text{vol}_{\partial U}^{n-1}$ 

where  $U^n$  is a compact region with boundary on a Riemannian  $M^n$ ,  $\text{vol}_{\partial U}^{n-1} = i^* \hat{N} \text{vol}^n$ , and  $\vec{v}$  is the velocity of the flow of  $U^n$ .

(Note: This is true even if  $M^n$  is not orientable since  $\text{vol}^n$  is a pseudo form and  $\vec{v}$  is a true vector)

(14) Let  $\vec{A}$  and  $\vec{B}$  be time dependant vector fields on  $\mathbb{R}^3$  and  $\rho(t, \vec{x})$  be continuous. Show that

$$(a) \frac{d}{dt} \int_{\mathbb{C}} \vec{A} \cdot d\vec{r} = \int_{\mathbb{C}} \left[ \frac{\partial \vec{A}}{\partial t} - \vec{v} \times \text{curl} \vec{A} + \text{grad}(\vec{v} \cdot \vec{A}) \right] \cdot d\vec{r}$$

(6)

$$(b) \frac{d}{dt} \int_S \vec{B} \cdot d\vec{S} = \int_S \left[ \frac{d\vec{B}}{dt} + (\nabla \cdot \vec{B})\vec{r} - \nabla \times (\vec{r} \times \vec{B}) \right] \cdot d\vec{S}$$

$$(c) \frac{d}{dt} \int_V \rho dV = \int_V \left[ \frac{d\rho}{dt} + \nabla \cdot (\rho\vec{r}) \right] dV$$

(15) Calculate for Riemannian manifold  $M^n$  with curvilinear coordinates (metric  $g_{ij}$ ):

$$(a) \Delta f = \nabla^2 f = \vec{\nabla} \cdot (\vec{\nabla} f)$$

$$(b) \nabla \times \vec{A}$$

using exterior methods

(16) Derive the following using vector identities

$$(a) \vec{\nabla}(fg) = f\vec{\nabla}g + g\vec{\nabla}f$$

$$(b) \nabla \cdot (f\vec{A}) = f\nabla \cdot \vec{A} + \vec{A} \cdot \vec{\nabla}f$$

$$(c) \nabla \times (f\vec{A}) = f\nabla \times \vec{A} + (\vec{\nabla}f) \times \vec{A}$$

$$(d) (\vec{A} \times \vec{B}) \cdot (\vec{C} \times \vec{D}) = ???$$

(17) Show that

$$\mathcal{L}_X \circ i_Y - i_Y \circ \mathcal{L}_X = i_{[X, Y]}$$